

Hilbert Space for Quantum Mechanics The Remaining Big Theorems

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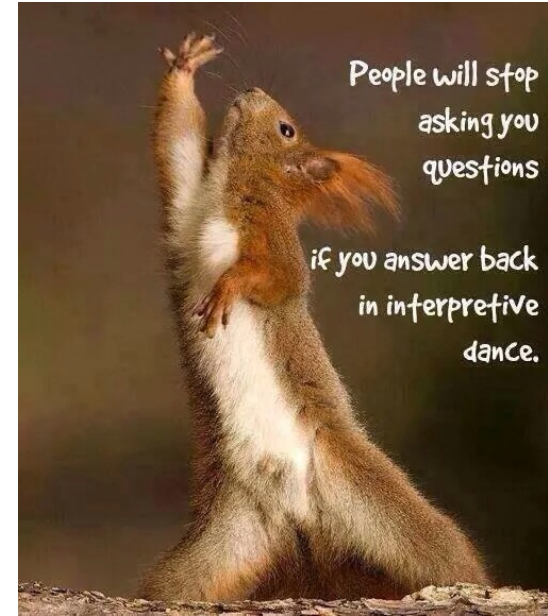


Table of Contents	Spectral Theorem:
Spectral Theorem: The	Projection-Valued Measure
Multiplication Operator Form	Form
Spectral Theorem: The	
Functional Calculus Form	Stones Theorem

MULTIPLICATION OPERATOR I

If μ is a finite measure on measure space M and $f: M \rightarrow \mathbb{R}$ is measurable define

$$\mathcal{D}_T = \{ g \in \mathcal{L}^2(M, \mu) \mid fg \in \mathcal{L}^2(M, \mu) \}$$

and $T: \mathcal{D}_T \rightarrow \mathcal{L}^2(M, \mu)$ by $T(g) = fg$.

Then T is self-adjoint and $\sigma(T)$ is the essential range of f ,

$$\sigma(T) = \{ \lambda \in \mathbb{R} \mid \mu(\{ m \in M \mid \lambda - \varepsilon < f(m) < \lambda + \varepsilon \}) > 0 \}.$$

If there are various operators like this in play we will call this one T_f .

MULTIPLICATION OPERATOR II

Suppose S is any self-adjoint operator on Hilbert space \mathcal{H} . Then there is a set M with finite measure μ and an isometry $U: \mathcal{H} \rightarrow \mathcal{L}^2(M, \mu)$ onto its range and a measurable function $f: M \rightarrow \mathbb{R}$ so that

- (i) $\psi \in \mathcal{D}_S$ if and only if $f U(\psi) \in \mathcal{L}^2(M, \mu)$.
- (ii) If $g \in U(\mathcal{D}_S)$ then $(USU^{-1}g)(m) = f(m)g(m)$ for all $m \in M$.
(a.e. wrt μ)

FUNCTIONAL CALCULUS II

Given all these properties it makes sense to denote $\phi(h)$ by $h(S)$.

(i)-(iv) imply, additionally, that:

- ▶ (v) If $S\psi = \lambda\psi$ then $\phi(h)\psi = h(\lambda)\psi$.
- ▶ (vi) If $h \geq 0$ then $\phi(h) \geq 0$.

FUNCTIONAL CALCULUS

- ▶ If S is any self-adjoint operator on Hilbert space \mathcal{H} there is a unique map ϕ from B , the bounded Borel functions on \mathbb{R} to $\mathcal{CL}(\mathcal{H})$ with the following properties.
- ▶ (i) ϕ is a $*$ -algebra homomorphism.
- ▶ (ii) ϕ is norm-continuous. In fact, $\|\phi(h)\| \leq \|h\|_\infty \forall h \in B$.
- ▶ (iii) If $h_n \rightarrow x \in B$ then $\phi(h_n)(\psi) \rightarrow S\psi$ on \mathcal{D}_S .
- ▶ (iv) If $h_n(x) \rightarrow h(x)$ for each x and $\|h_n\|_\infty$ is bounded then $\phi(h_n) \rightarrow \phi(h)$ strongly.

PROJECTION-VALUED MEASURE I

If S is any self-adjoint operator on Hilbert space \mathcal{H} and χ_A is the characteristic function of Borel $A \subset \mathbb{R}$ define $P_A = \phi(\chi_A) = \chi_A(S)$. And define $P_t = P_{(-\infty, t]}$.

- ▶ Since $\chi_A^2 = \chi_A$ and χ_A is real valued each P_A is an orthogonal projection, $P_\emptyset = 0$ and $P_{\mathbb{R}} = I$, the identity on \mathcal{H} .
- ▶ $P_A P_B = P_{A \cap B}$.
- ▶ (i) If $A = \bigcup_{i=1}^\infty A_i$ and $A_i \cap A_j = \emptyset$ whenever $i \neq j$ then P_A is the strong limit of $\sum_{i=1}^n P_{A_i}$.

PROJECTION-VALUED MEASURE II

- ▶ For each fixed ψ the inner product $\langle \psi, P_A(\psi) \rangle$ is an ordinary measure on \mathbb{R} and we use $d\langle \psi, P_A(\psi) \rangle$ to denote integration with respect to that measure.
- ▶ For each ψ the map

$$f \rightarrow \int f(t) d\langle \psi, P_t(\psi) \rangle$$

is bounded and linear, and so corresponds to inner product against some $\eta = B\psi$, as in $\langle \psi, B\psi \rangle = \int f(t) d\langle \psi, P_t(\psi) \rangle$.

- ▶ We denote this B by $B = \int f(t) dP_t$ and show that $S = \int t dP_t$.

PROJECTION-VALUED MEASURE III

Theorem

There is a one-to-one correspondence between self-adjoint operators T and projection-valued measures P given by

$$S \rightarrow P$$

$$P \rightarrow \int t dP_t.$$

STONE'S THEOREM I

- ▶ A strongly continuous one-parameter group of unitary operators on Hilbert \mathcal{H} is a map $U: \mathbb{R} \rightarrow \mathcal{C}\mathcal{L}(\mathcal{H})$ which satisfies
 - $U(t)$ is unitary and $U(s+t) = U(s)U(t)$ for $s, t \in \mathbb{R}$.
 - If $t_i \rightarrow a$ then $U(t_i) \rightarrow U(a)$ strongly.
- ▶ If T is self-adjoint define $U(t) = e^{itT}$. Then U is a strongly continuous one-parameter group of unitary operators and, in addition,
 - For each $\psi \in \mathcal{D}_T$ we have $\frac{U(t)\psi - \psi}{t} \rightarrow iT\psi$ as $t \rightarrow 0$ and
 - If $\lim_{t \rightarrow 0} \frac{U(t)\psi - \psi}{t}$ exists then $\psi \in \mathcal{D}_T$.

STONE'S THEOREM II

Theorem

If U is any strongly continuous one-parameter group of unitary operators then there is a self-adjoint operator T with $U(t) = e^{itT}$.

T is called the infinitesimal generator of U .

The phrasing of these results, and their order, was adapted from the superb standard reference by Reed and Simon, *Functional Analysis*, Academic Press, 1980. I **highly** recommend it as a very clean mathematical presentation with pedagogical flair, filled with insight into the sticky issues involved in applications to quantum physics.